



# Quantum control of interacting qubits

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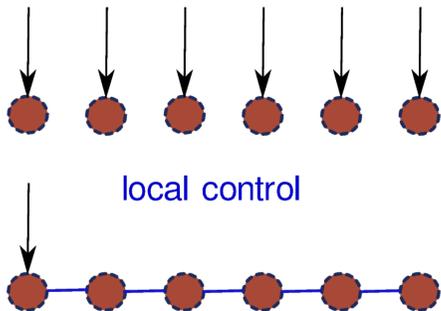
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## Local control approach for interacting systems

### "conventional" control



### Complete controllability:

Acting locally on one of the end spins of an  $XXZ$ -Heisenberg spin chain ensures complete controllability.

[see, e.g., D. Burgarth *et al.*, PRA **79**, 060305(R) (2009)]

## $XXZ$ Heisenberg spin-1/2 chain

Total Hamiltonian:  $H(t) = H_0 + H_c(t)$

$XXZ$  Heisenberg coupling with anisotropy  $\Delta$  (chain length  $N_s$ ):

$$H_0 = J \sum_{i=1}^{N_s-1} (S_{ix}S_{i+1,x} + S_{iy}S_{i+1,y} + \Delta S_{iz}S_{i+1,z})$$

Zeeman-like control part (control fields acting on the first spin only):

$$H_c(t) = h_x(t)S_{1x} + h_y(t)S_{1y}$$

Dynamical Lie algebra generated by  $iH_0, iS_{1x}, iS_{1y}$  has dimension  $d^2 - 1$  ( $d = 2^{N_s}$ )  $\Rightarrow$  **complete controllability**

## Operator control

Realize a desired unitary transformation (target quantum gate)  $U_{\text{target}}$  at a time  $t = t_f$  (more general than state control).

## Control objectives (target gates)

One-qubit gate:

$X$  (flip on the last qubit)

Entangling two-qubit gates:

**CNOT** and  $\sqrt{\text{SWAP}}$  (on the last two qubits)

## Gate fidelity is defined as

$$F(t_f) = \frac{1}{d} \left| \text{tr}[U^\dagger(t_f)U_{\text{target}}] \right|$$

$U(t_f)$ : time-evolution operator at  $t = t_f$

## THREE-SPIN SYSTEM

## Control fields implementing a flip of the third qubit

Piecewise-constant control fields:  $N_t$  pulses of length  $T$  (total evolution time  $t_f = N_t T$ )

Maximization of fidelity with respect to the  $N_t$  field amplitudes  $\Rightarrow$  **optimal control sequences** (gate error:  $< 10^{-3}$ )

**Estimate of minimal gate times**

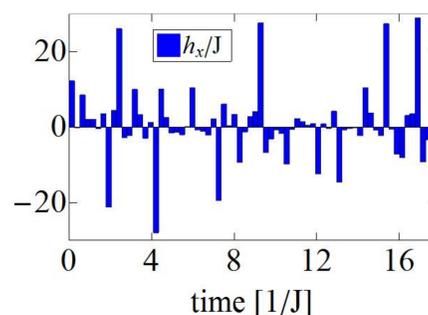
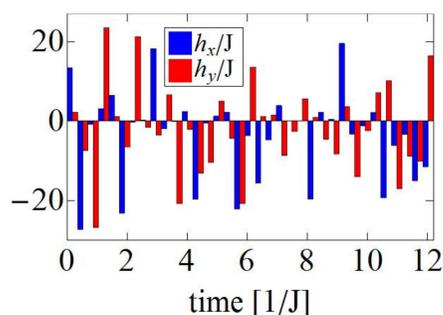
### Alternate x and y pulses:

$h_{x,j}, h_{y,j}$  ( $j = 1, \dots, N_t/2$ )

Lie completion:  $[S_{1x}, S_{1y}] = iS_{1z}$

$\Rightarrow$

System is completely controllable, i.e., the dimension of the dynamical Lie algebra is 63.



### Controls in x direction only:

$h_{x,j}$  ( $j = 1, \dots, N_t$ )

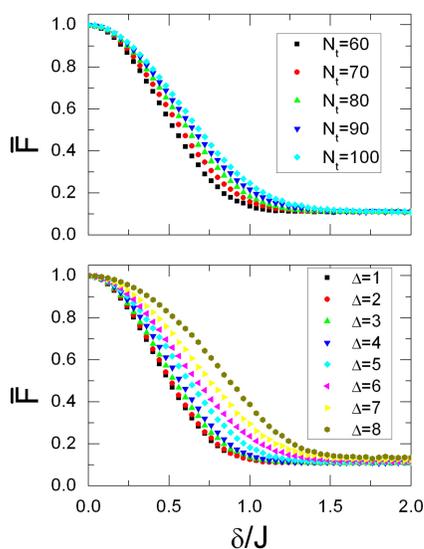
Dimension of the dynamical Lie algebra: 30 for  $\Delta \neq 1$  and 18 for  $\Delta = 1$

Only specific gates can be realized, such as  $X$  and  $\sqrt{\text{SWAP}}$  gates.

## Robustness of fidelity to random errors in the control fields

## Spectral filtering (smoothing) of control fields

$\bar{F}$ : fidelity averaged over random realizations  
 $\delta$ : half-width of uniform random distribution  
 Gate implementation using **x and y controls**



$\sqrt{\text{SWAP}}$  on the second and third qubit ( $\Delta = 1.2, t_f = 60J^{-1}$ )

**Fixed  $t_f$ : Fidelity less sensitive to random errors for faster switching (larger  $N_t$ ).**

**CNOT** on the second and third qubit ( $N_t = 70, T = 1J^{-1}$ )

**Robustness depending on  $\Delta$ : Larger  $\Delta$  yield higher fidelity.**

**Universal saturation at a value of  $1/d = 0.125$  for any  $t_f$  (in the completely-controllable case)**

[see R. Heule *et al.*, PRA **82**, 052333 (2010), and arXiv:1010.5715 for more details]

Frequency-filtered control fields:

$$\tilde{h}_j(t) = \mathcal{F}^{-1}[f(\omega)\mathcal{F}[h_j(t)]] \quad (j = x, y)$$

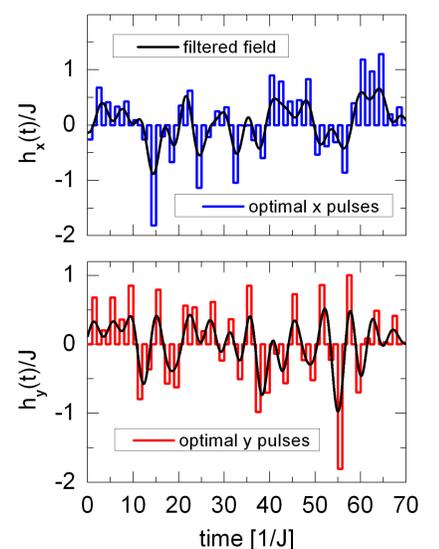
Ideal low-pass filter:

$$f(\omega) = \theta(\omega + \omega_0) - \theta(\omega - \omega_0)$$

Example: implementation of **CNOT** on the second and third qubit with a gate error of  $10^{-8}$

**Filtered control field retains a fidelity of 0.9.**

(with a cut-off  $\omega_0/J = \pi/2$ )



## Conclusions and Outlook

- Minimal gate times strongly depend on  $\Delta$  and grow rapidly with the chain length.
- Relatively high gate fidelities possible in the presence of random errors and for smoothed control fields.
- Generalizations: spins  $s > 1/2$ , more complex control pulses, open-system scenario.