

Anomalous suppression of the shot noise in a nanoelectromechanical system

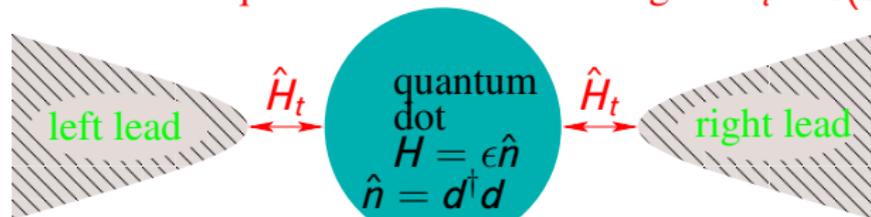
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08.08.2006

cond-mat/0607080

The system and its Hamiltonian

sequential electron tunneling $\hat{H}_t = t(c^\dagger d + d^\dagger c)$



$$\hat{H}_{n,b} = \lambda \omega_0 (b^\dagger + b) \hat{n}$$

Oscillator
(vibrating device)

$$H_b = \omega_0 (b^\dagger b + 1/2)$$

$$\hat{H}_{b,env} = \sum_j \chi_j \omega_j (a_j^\dagger + a_j) (b^\dagger + b)$$

Thermal bath

$$H_{env} = \sum_j \omega_j (a_j^\dagger a_j + 1/2)$$

characterized
by the
"spectral function"

$$J(\omega) = 2\pi \sum_j \omega_j^2 \chi_j^2 \delta(\omega - \omega_j)$$

Task: calculate Fano factor: $F = \frac{S(\omega=0)}{2e\langle I \rangle}$.

Rate equation

occupation numbers: $|n, l\rangle$, n = electrons in the dot,
 l = phonons in the oscillator.

Considering “adiabatic” regime:

$\omega_0 \gg \Gamma^{(0)} = 2\pi\nu t_0^2$ = tunneling rate.

\Rightarrow ignoring off-diagonal elements of the density matrix:

$\langle n, l | \hat{\rho} | n', l' \rangle \propto \delta_{n,n'} \delta_{l,l'}$.

$i\frac{\partial \hat{\rho}}{\partial t} = [\hat{H}, \hat{\rho}] \Rightarrow$ equation of motion:

$$\frac{d}{dt} P_{nl} = \sum_{n' \neq n} \sum_{l', \alpha} [P_{n'l'} \Gamma_{\alpha l' \rightarrow l}^{n' \rightarrow n} - P_{nl} \Gamma_{\alpha l \rightarrow l'}^{n \rightarrow n'}] + \sum_{l'} [P_{nl'} \Gamma_{l' \rightarrow l}^{rel} - P_{nl} \Gamma_{l \rightarrow l'}^{rel}].$$

← relaxation rates $\Gamma_{l' \rightarrow l}^{rel}$ are considered separately.

← reminds quasiclassical kinetic equation for the distribution function.

Now let us evaluate the “collision term” on the rhs
(i.e., the rates Γ)

Expressions for the rates Γ

The transition rates are evaluated using **Fermi golden rule approximation**. Now let us evaluate the “collision term” on the rhs (i.e., the rates Γ).

- ▶ Relaxation rates:

$$\Gamma_{I \rightarrow (I-1)}^{rel} = e^{\beta\omega_0} \Gamma_{(I-1) \rightarrow I}^{rel} = I \frac{\mathcal{J}(\omega_0)}{1 - e^{-\beta\omega_0}},$$



–equilibrium expression!

- ▶ Charge transfer rates:

$$\begin{aligned}\Gamma_{\alpha I \rightarrow I'}^{0 \rightarrow 1} &= \Gamma^{(0)} X_{I'I} f_{\alpha}(\omega_0(I' - I)), \\ \Gamma_{\alpha I \rightarrow I'}^{1 \rightarrow 0} &= \Gamma^{(0)} X_{I'I} [1 - f_{\alpha}(\omega_0(I - I'))], \\ X_{I'I} &= e^{-\lambda^2} \lambda^{2|I-I'|} \frac{I_{<}!}{I_{>}!} |L_{I_{<}}^{|I-I'|}(\lambda^2)|^2,\end{aligned}$$

Using these (known before) expressions, the authors numerically solve the rate equation.

Limiting cases

- ▶ $\lambda \rightarrow 0$ sequential tunneling, minimal Fano-factor = $1/2$.
- ▶ $\lambda \gg 1$

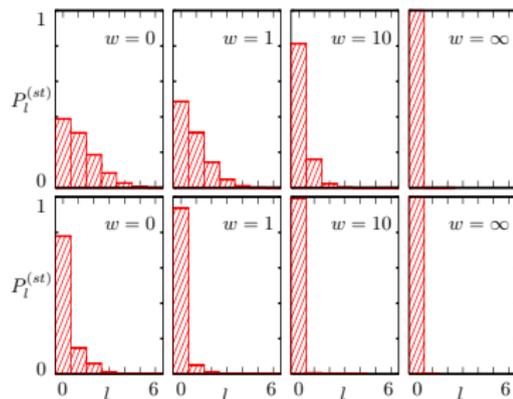


Figure: Stationary phonon probability distribution $P_l^{(st)}$ for different values of w . Upper panels: $\lambda^2 = 0.4$, lower: $\lambda^2 = 7$.

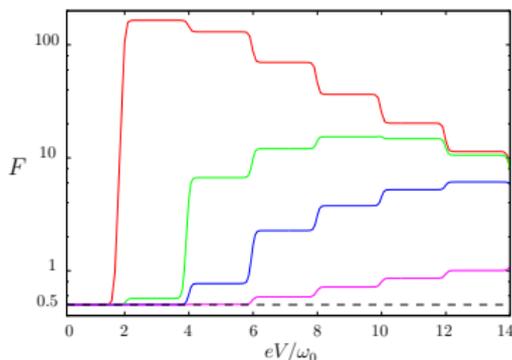


Figure: Fano factor as a function of voltage for $\lambda^2 = 16$ and for different values of the relaxation strength w ; red $w = 0$, green $w = 0.1$, blue $w = 1$, magenta $w = 10$.

What happens for $\lambda^2 \sim 1$?

Interpolation between two limiting cases?

No! Finite relaxation rate can induce unexpected features!

Moderate coupling

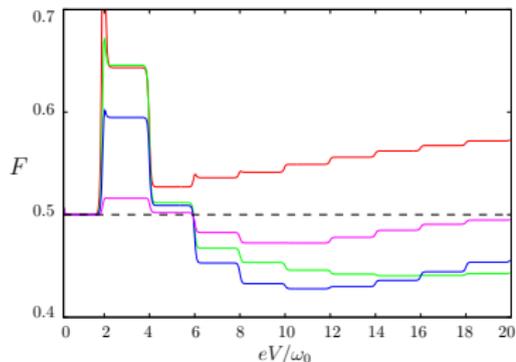


Figure: Fano factor as a function of voltage at $\lambda^2 = 3$ and for different values of the relaxation strength w ; red $w = 0$, green $w = 5$, blue $w = 12$, magenta $w = 100$. Dashed line, $F = 1/2$. Other parameters: $\bar{\epsilon} = 0$ and $k_B T = 0.02 \omega_0$.

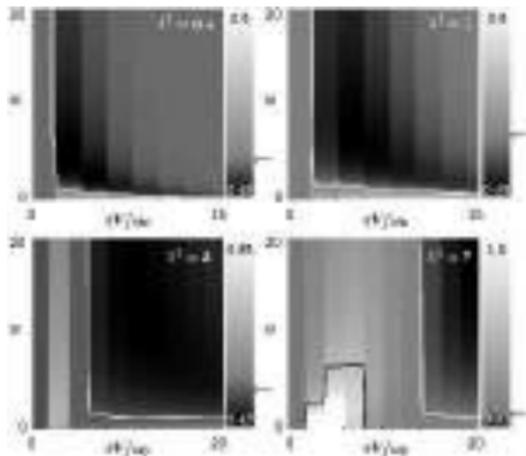


Figure: Fano factor as a function of bias V and relaxation w for different values of λ . In all the panels: dark gray $F < 1/2$, medium gray $F = 1/2$ (indicated by the arrow in the color map) and light gray $F > 1/2$. The white line corresponds to $F = 1/2$. The black line in the 4th panel delimits the region where noise is superpoissonian, $F > 1$. Other parameters: $\bar{\epsilon} = 0$ and $k_B T = 0.02 \omega_0$.