

Theoretical Solid-State Physics, Herbstsemester 2012

Blatt 6

Abgabe: 01.11.2012, 12:00H (Treppenhaus 4. Stock)

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(1) **Noninteracting Green's function for topological insulators** (6 Points)

Recall the derivation of the noninteracting Green's function for a system of fermions. Consider now two-dimensional surface electrons (in the $x - y$ plane) described by the Hamiltonian

$$H_{\text{TI}} = \hbar v_F \sum_{\mathbf{k}, \alpha, \beta} c_{\alpha\mathbf{k}}^\dagger \mathbf{k} \cdot \boldsymbol{\sigma}_{\alpha\beta} c_{\beta\mathbf{k}}, \quad (1)$$

where the indices α and β represent the electron spin (\uparrow or \downarrow), $\boldsymbol{\sigma} \equiv (\sigma_x, \sigma_y)$ is a vector of Pauli spin matrices, $\mathbf{k} \equiv (k_x, k_y)$ is the electron wave vector, and v_F the Fermi velocity. In this problem, we are interested in calculating the Green's function in the ground state ($T = 0$).

- (a) Using the equations of motion for operators in the Heisenberg picture, determine the time-dependence of $c_{\sigma\mathbf{k}}^\dagger(t)$ and $c_{\sigma\mathbf{k}}(t')$ ($\sigma \in \{\uparrow, \downarrow\}$) for the noninteracting system described by H_{TI} .
- (b) Using part (a), show that the noninteracting retarded Green's function

$$G_{\alpha\beta}^R(\mathbf{k}, t - t') \equiv -i\Theta(t - t') \langle [c_{\alpha\mathbf{k}}^\dagger(t), c_{\beta\mathbf{k}}(t')] \rangle, \quad (2)$$

which in this case is a 2×2 matrix in spin space, at $T = 0$ is given by

$$G^R(\mathbf{k}, t - t') = -i\Theta(t - t') e^{-iv_F \mathbf{k} \cdot \boldsymbol{\sigma} (t - t')}. \quad (3)$$

- (c) Show that in the frequency domain

$$\begin{aligned} G^R(\mathbf{k}, \omega) &\equiv \int d(t - t') e^{i(\omega + i\eta)(t - t')} G^R(\mathbf{k}, t - t') \\ &= \frac{v_F \mathbf{k} \cdot \boldsymbol{\sigma} - \omega - i\eta}{v_F^2 k^2 - \omega^2}, \end{aligned} \quad (4)$$

where $\eta \rightarrow 0^+$.

(2) **Momentum dependence of electron-phonon coupling** (4 Points)

Given below are three different types of short-range coupling of a single electron with Einstein phonons (or, more generally, dispersionless bosons) of frequency ω , on a one-dimensional lattice (lattice constant $\equiv 1$) with N sites. The coupling constants g , ϕ , and ϕ_b are dimensionless. For each of these couplings, find the equivalent momentum-space

representation, i.e., transform the corresponding electron-phonon coupling Hamiltonian in real space to the form

$$H_{\text{e-ph}} = \frac{1}{\sqrt{N}} \sum_{k,q} \gamma(k, q) a_{k+q}^\dagger a_k (b_{-q}^\dagger + b_q). \quad (5)$$

Here the a 's and b 's are the electron- and phonon operators, respectively, while k and q are the corresponding quasimomenta.

(a) Holstein-type (purely local) coupling:

$$H_{\text{e-ph}} = g\omega \sum_{i=1}^N a_i^\dagger a_i (b_i^\dagger + b_i) \quad (6)$$

(b) Su-Schrieffer-Heeger (Peierls-type) coupling:

$$H_{\text{e-ph}} = \phi\omega \sum_{i=1}^N (a_{i+1}^\dagger a_i + \text{h.c.}) (b_{i+1}^\dagger + b_{i+1} - b_i^\dagger - b_i) \quad (7)$$

(c) “breathing” coupling (relevant in cuprate high- T_c superconductors):

$$H_{\text{e-ph}} = \phi_b\omega \sum_{i=1}^N a_i^\dagger a_i (b_{i-1/2}^\dagger + b_{i-1/2} - b_{i+1/2}^\dagger - b_{i+1/2}). \quad (8)$$

Here, $i \pm 1/2$ refers to the fact that the Einstein oscillators are placed in the middle between two sites.

Then comment on the differences between these three types of electron-phonon interaction as far as the momentum dependence of the vertex function $\gamma(k, q)$ is concerned.